

The magnetic field strength of neutron stars

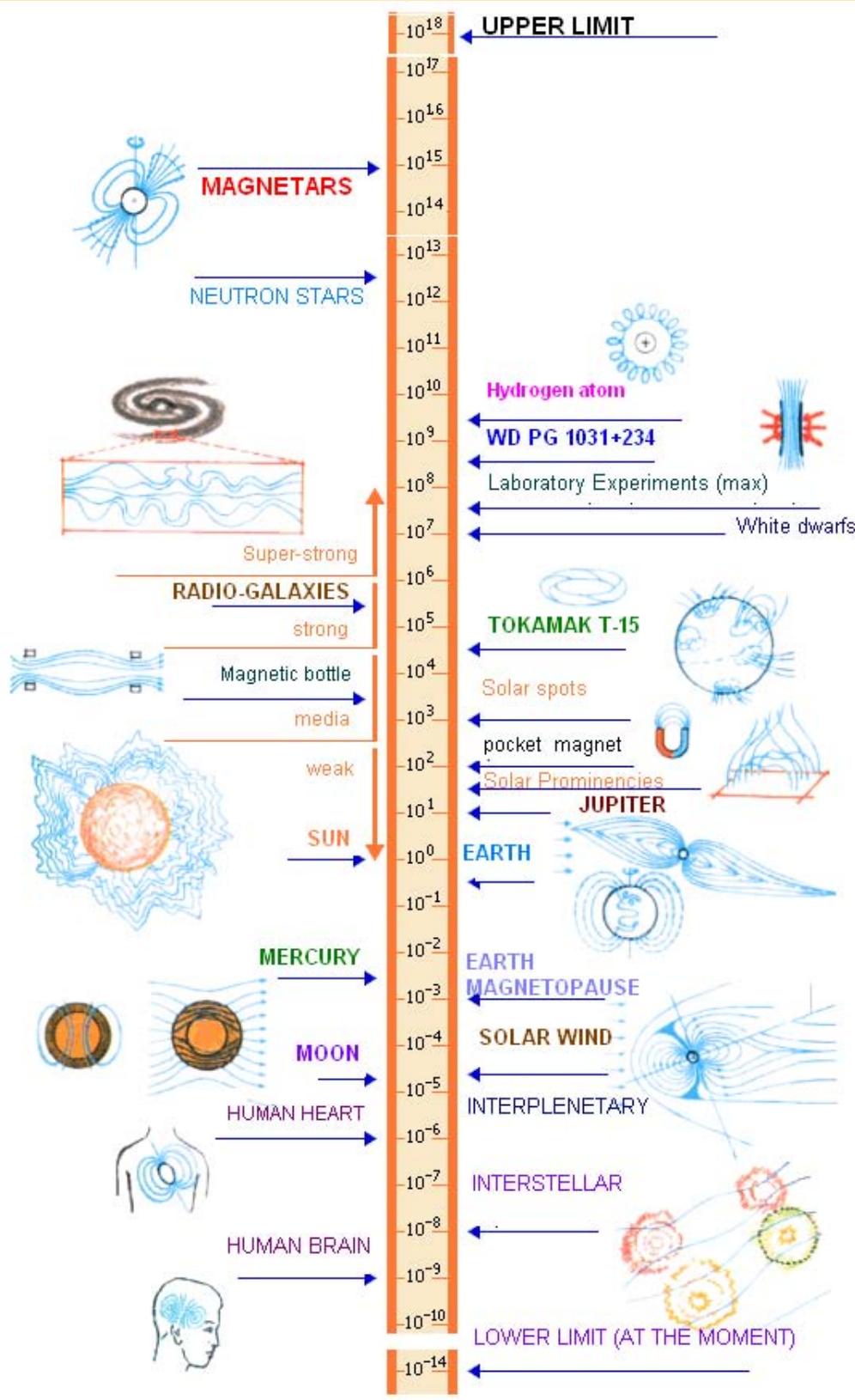
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2. Ikhsanov N.R. & Beskrovnaya, N.G. 2012 Astronomy Reports, **56**, 589
3. Ikhsanov N.R. & Finger, M.H. 2012 ApJ, **753**, 1
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UPPER LIMIT

Gravitational Energy \equiv Magnetic energy

$$B_{\max} \sim 10^{18} \text{ G}$$

QUANTUM CRITICAL FILED

$$B_{\text{cr}} = \frac{m_e^2 c^3}{e \hbar} \simeq 4.4 \times 10^{13} \text{ G}$$

REGULAR FIELD OF NEUTRON STARS

Magnetic Flux Conservation during Collapse

$$B_{\text{ns}} \sim B_{\text{in}} \left(\frac{R_{\text{in}}}{R_{\text{ns}}} \right)^2 \simeq 10^{12} \text{ G} \times$$

$$\times \left(\frac{B_{\text{in}}}{100 \text{ G}} \right) \left(\frac{R_{\text{in}}}{10^6 \text{ km}} \right)^2 \left(\frac{R_{\text{ns}}}{10 \text{ km}} \right)^{-2}$$

Magnetic Field Measurements

There are four ways of constraining the pulsars B field

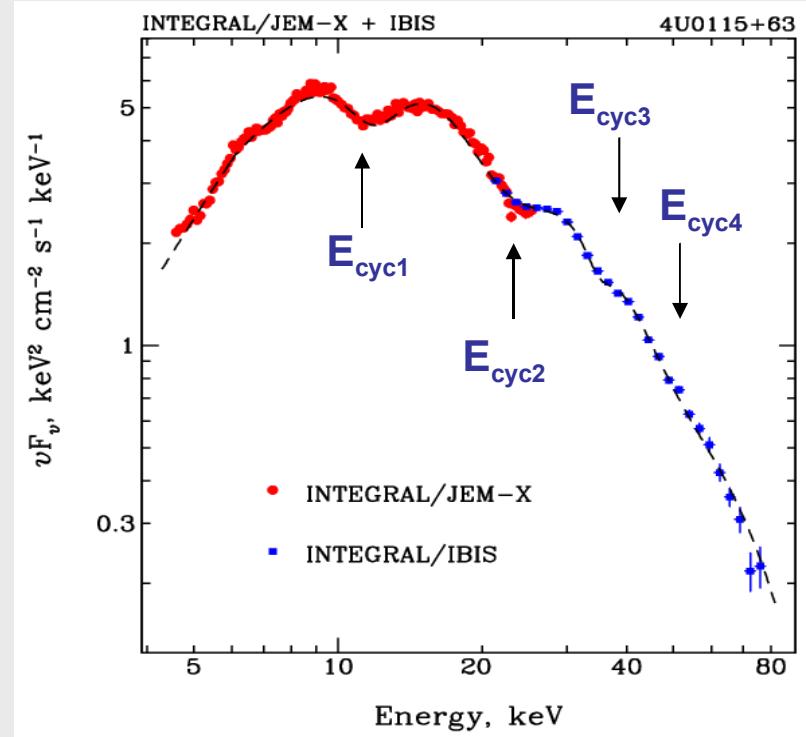
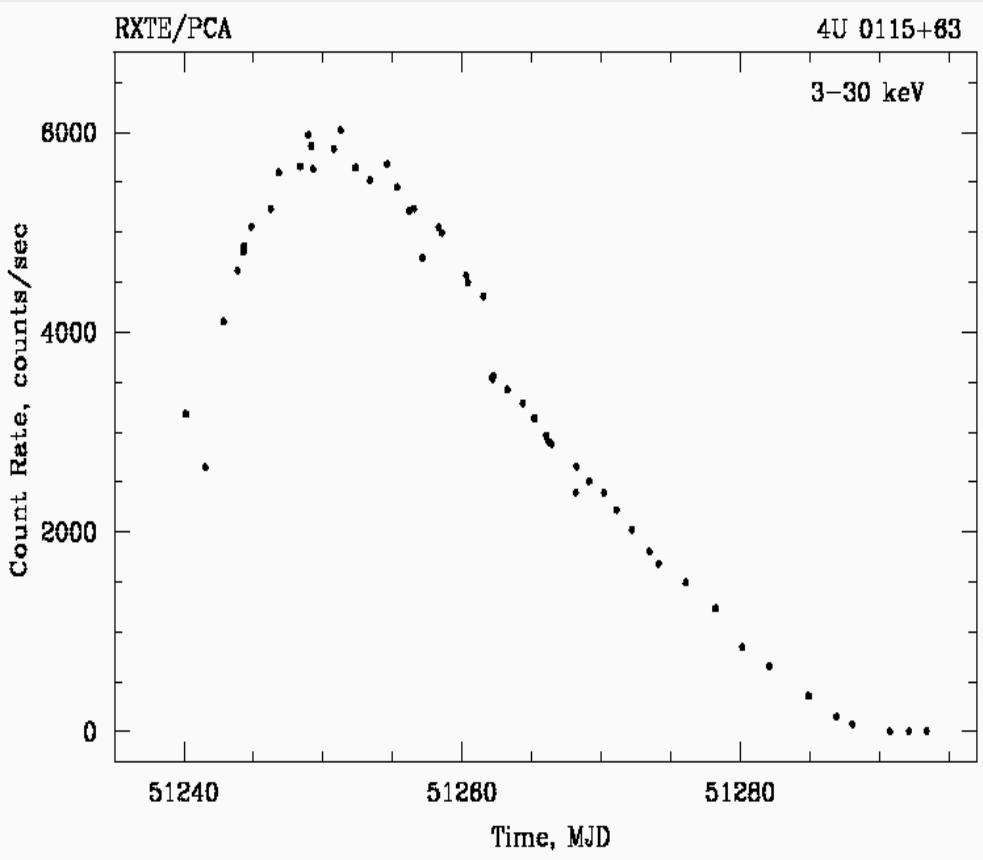
- Measure the lowest cyclotron absorption line energy

$$B_{\text{CRSF}} \leq 10^{12} \left(\frac{E_{\text{cyc}}}{11.6 \text{ keV}} \right) \text{ G}$$

- Measure the spin-down rate of pulsars;
- Measure the spin-up rate at high luminosity of the accretion-powered X-ray pulsars;
- Observations of the cutoff of accretion due to centrifugal inhibition;

Вспышки излучения 4U 0115+63 1999 и 2004 гг.

Цыганков и др. (2007)



Параметры спектра

$$\alpha = 0.09 \pm 0.01$$

$$E_{\text{cut}} = 8.9 \pm 0.1 \text{ кэВ}$$

$$E_{\text{cyc}1} = 11.16 \pm 0.03 \text{ кэВ}$$

$$E_{\text{cyc}2} = 21.2 \pm 0.1 \text{ кэВ}$$

$$E_{\text{cyc}3} = 34.6 \pm 0.2 \text{ кэВ}$$

$$E_{\text{cyc}4} = 44.9 \pm 0.3 \text{ кэВ}$$

$$E_{\text{cyc}1}/E_{\text{cyc}2}/E_{\text{cyc}3}/E_{\text{cyc}4} = 1/1.9/3.1/4.0$$

Максимальный диапазон светимостей
 $\sim 1 \times 10^{37} - \sim 2 \times 10^{38}$ эрг/с
(d~7 кпс)

Зависимость доли пульсирующего излучения от светимости и энергетического диапазона

*Временной анализ данных обсерватории ИНТЕГРАЛ/IBIS в узких
энергетических каналах.*

Только яркие источники (поток в 20-50 кэВ > ~100 mCrab):

Источник

4U 0115+63

V 0332+53

A 0535+262

Hercules X-1

Vela X-1

GX 301-2

GX 1+4

Cen X-3

AO 1657-415

EXO 2030+375

E_{cyc}

~11, 21, 34, 45 кэВ

~28, 55 кэВ

~45, 100 кэВ

~38 кэВ

~22, 56 кэВ

~50 кэВ

-

~31 кэВ

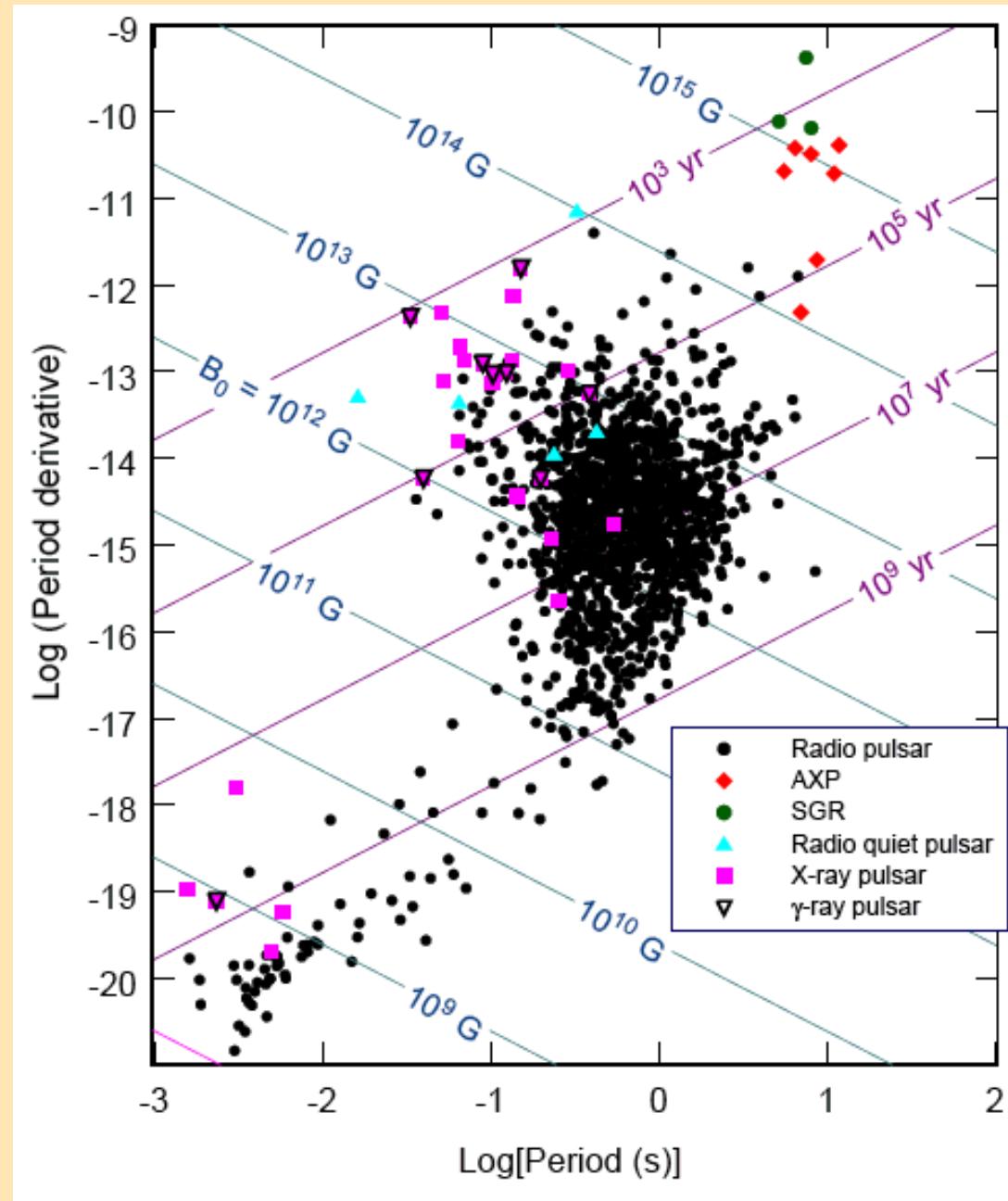
~78 кэВ

-

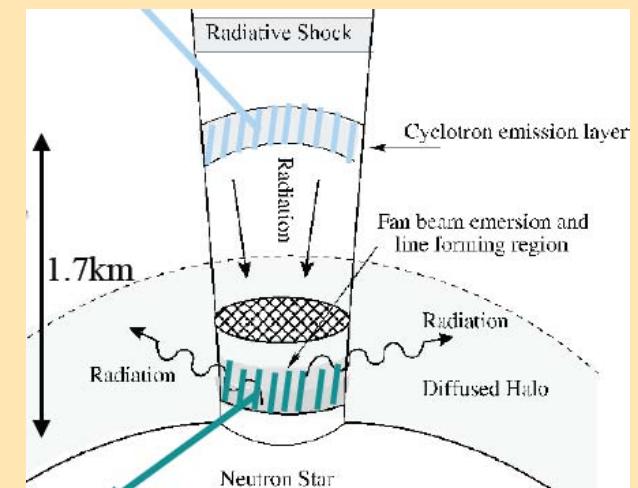
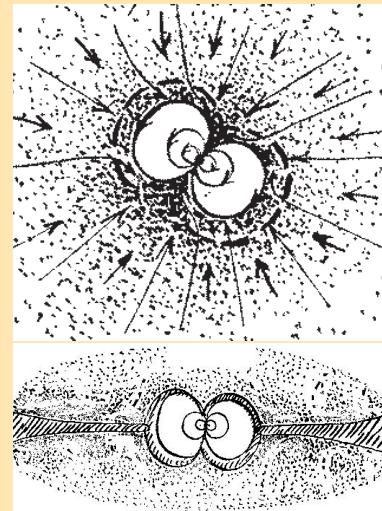
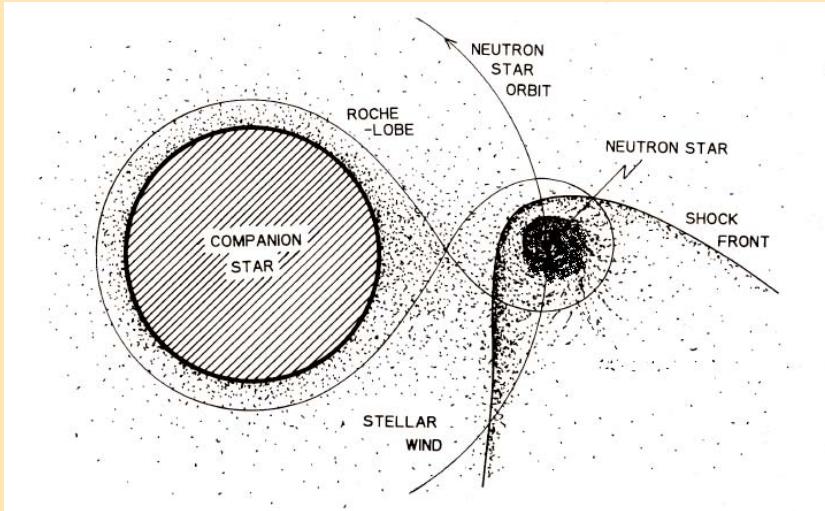
Cyclotron Line Measurements

Class	Name	E_{cyc} (keV)	Reference
Be/X-ray	EXO 2030+375	10.5	Wilson et al. 2007
	4U 0115+63	11.5	White, Swank & Holt, 1983
	V 0332+53	28.5	Makishima et al., 1991
	X Per	28.6	Colburn et al., 2001
	Cep X-4	30.0	Mihara et al. 1991
	XTE J1946+274	34.9	Heindl et al. 2001
	MX 0656-072	36.0	McBride, V.A. et al., 2005
	A 0535+26	47.0	Caballero et al 2007
Disk Fed	Cen X-3	30.0	Burderi et al., 2000
	4U 1626-67	39.0	Coburn et al. 2000
	Her X-1	42.0	Manchanda, 2003
Wind Fed	4U 1907+09	18.3	Cusumano et al. 1998
	4U 1538-52	20.7	Robba et al. 2001
	2S 0114+65	22.0	Bonning et al. 2005
	Vela X-1	23.3	Kreykenbohm et al., 2002
	4U 2206+54	32.0	Blay et al., 2006
	GX 301-2	42.4	La Barbera et al. 2005

Magneto-dipole waves $L_{\text{md}} = f_{\text{m}} \frac{\mu^2 \omega_s^4}{c^3} = I \omega \dot{\omega}$



Accretion-powered pulsars in wind-fed HMXBs



MASS CAPTURE

Relative velocity

$$V_{\text{rel}} = c_s + |\vec{V}_{\text{orb}} + \vec{V}_w|$$

Bondi radius

$$r_G = \frac{2GM_{\text{ns}}}{V_{\text{rel}}^2}$$

Mass capture rate

$$\dot{\mathfrak{m}}_c = \pi r_G^2 \rho_\infty V_{\text{rel}}$$

MASS ACCRETION

Mass accretion rate

$$\dot{\mathfrak{m}}_a = \frac{L_x R_{\text{ns}}}{GM_{\text{ns}}}$$

Magnetic field

$$B_{\text{CRSF}} \leq 10^{12} \left(\frac{E_{\text{cyc}}}{11.6 \text{ keV}} \right) \text{ G}$$

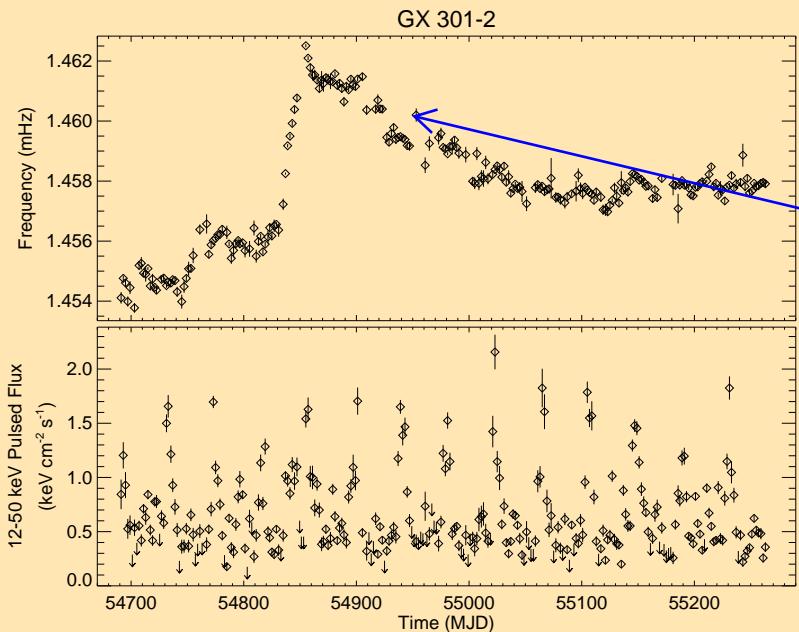
Dipole magnetic moment

$$\mu \simeq \frac{1}{2} B_{\text{CRSF}} R_{\text{ns}}^3$$

- Accretion of mass

- Accretion of angular momentum

Spin Evolution (Spin-up \longleftrightarrow Spin-down) $\nu = 1/P_s$



$$2\pi I \dot{\nu} = K_{su} - K_{sd}$$

$$|K_{sd}| \geq 2\pi I |\dot{\nu}_{sd}|$$

$$|K_{sd}(r_A)| \leq \begin{cases} |K_{sd}^{(0)}| = \mu^2/r_{\text{cor}}^3 \\ |K_{sd}^{(t)}| = \mathfrak{M} \omega_s r_A^2 \end{cases}$$

Name	B_{12} , G	P_s , s	$ \dot{\nu}_{sd}^{\text{obs}} $	$ \dot{\nu}_{sd}^{(0)} $	$ \dot{\nu}_{sd}^{(t)} $
OAO 1657-415	3.2	38	3.2	0.056	0.18
Vela X-1	2.6	283	0.2	0.0003	0.012
4U 1907+09	2.1	438	0.04	0.0002	0.008
4U 1538-522	1.8	529	0.06	0.0002	0.008
GX 301-2	4.0	683	0.1	0.003	0.02

The Problem (since 1973)

The **spin-down torque** applied to Long-Period X-ray Pulsars in HMXBs

is a factor of 10 larger

than the **maximum possible spin-down torque** expected within
the **conventional scenario** of accretion from a **non-magnetized stellar wind**.

Conventional scenario:

- Magnetospheric radius: $r_m \equiv r_A = \left(\frac{\mu^2}{\dot{M} \sqrt{GM_{ns}}} \right)^{2/7}; \quad \left[\dot{M} = \frac{L_X R_{ns}}{GM_{ns}} \right]$

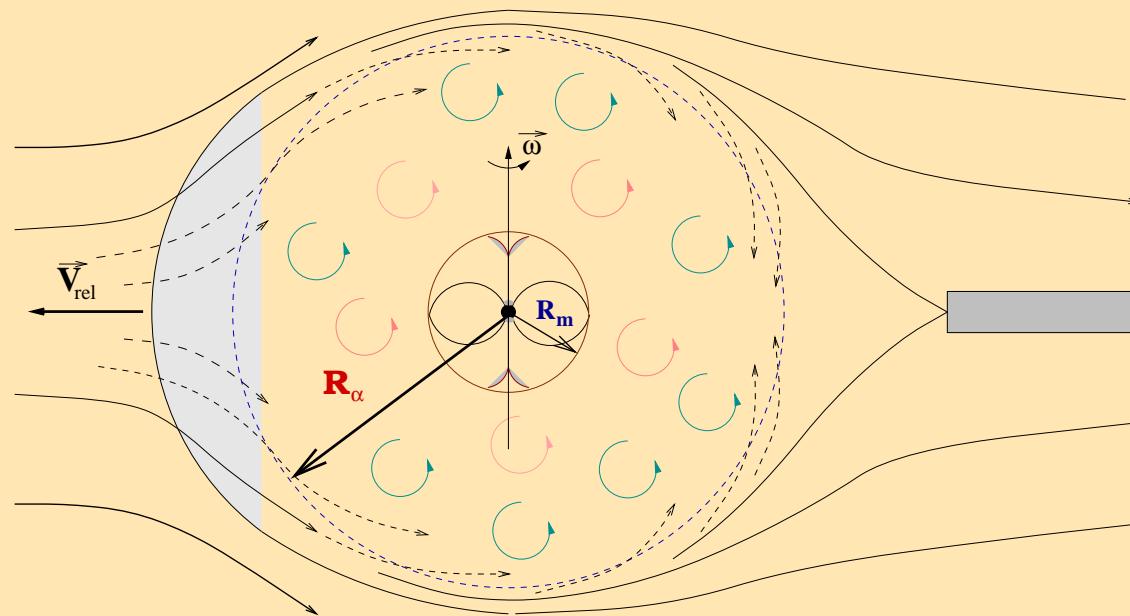
- Spin-down torque: $|K_{sd}(r_A)| \sim \begin{cases} |K_{sd}^{(0)}| = \mu^2 \times \left[\frac{4\pi^2 k_t}{GM_{ns} P_s^2} \right] \\ |K_{sd}^{(t)}| = \mu^{8/7} \times \left[\frac{2\pi k_t L_X^{3/7} R_{ns}^{3/7}}{(GM_{ns})^{5/7} P_s} \right] \end{cases}$

Torque definition

$$\mathbf{K} = \mathbf{F} \times \mathbf{r}$$

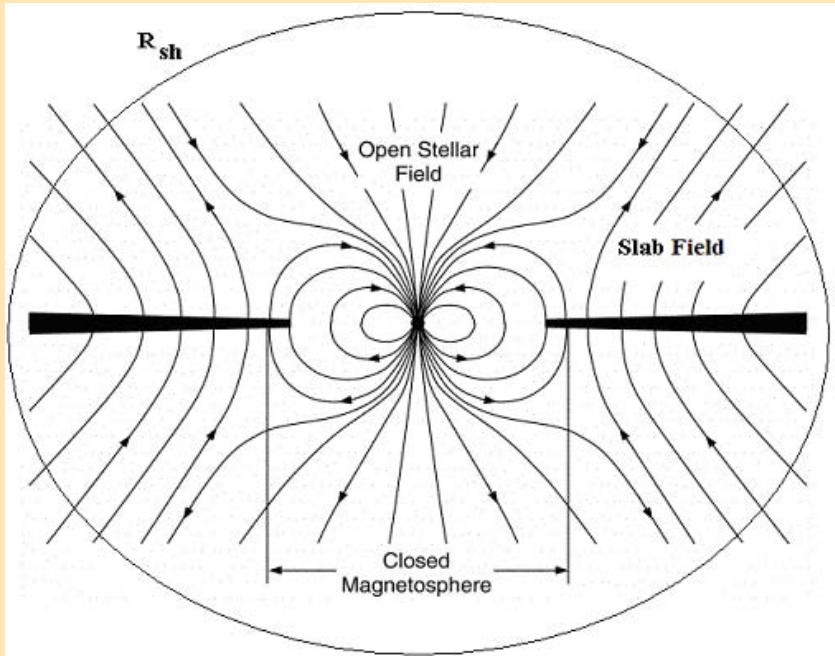
Model task:

A sphere of the radius r_m
is rotating in a viscose medium



Spin-down torque in general case (I of II steps)

$$|K_{\text{sd}}(r_m)| = \nu_t S_{\text{eff}}(r_m) \rho(r_m) v_\phi(r_m)$$



Viscosity

$$\nu_t = k_t \ell_t v_t$$

Effective Area

$$S_{\text{eff}} = 2\pi r_m h_s(r_m)$$

Thickness

$$h_s(r_m) = c_s^2(r_m) \left(\frac{r_m^2}{GM_{\text{ns}}} \right)$$

Density

$$\rho(r_m) = \frac{1}{c_s^2(r_m)} \left(\frac{\mu^2}{2\pi r_m^6} \right)$$

ϕ -velocity

$$v_\phi = r_m [\omega_s - \Omega(r_m)]$$

$$|K_{\text{sd}}(r_m)| = [k_t \ell_t v_t] \left[\frac{\mu^2}{r_m^2} \frac{\omega_s}{GM_{\text{ns}}} \right] \left(1 - \frac{\Omega(r_m)}{\omega_s} \right)$$

Spin-down torque in general case (II of II steps)

$$|K_{\text{sd}}(r_m)| = [k_t \ell_t v_t] \left[\frac{\mu^2}{r_m^2} \frac{\omega_s}{GM_{\text{ns}}} \right] \left(1 - \frac{\Omega(r_m)}{\omega_s} \right)$$

$$\ell_t \leq r_m$$

$$v_t \leq v_k(r_m)$$

$$r_{\text{cor}} = \left(\frac{GM_{\text{ns}}}{\omega_s^2} \right)^{1/3}$$

$$|K_{\text{sd}}(r_m)| = \left[k_t r_m \left(\frac{GM_{\text{ns}}}{r_m} \right)^{1/2} \right] \left[\frac{\mu^2}{r_m^2} \frac{\omega_s}{GM_{\text{ns}}} \right] \left(1 - \frac{\Omega(r_m)}{\omega_s} \right)$$

$$= k_t \frac{\mu^2}{r_m^{3/2}} \left[\frac{\omega_s}{(GM_{\text{ns}})^{1/2}} \right] \left(1 - \frac{\Omega(r_m)}{\omega_s} \right)$$

$$|K_{\text{sd}}(r_m)| = k_t \frac{\mu^2}{(r_m r_{\text{cor}})^{3/2}} \left(1 - \frac{\Omega(r_m)}{\omega_s} \right)$$

Ikhsanov N.R. & Beskrovnyaya, N.G. 2012 *Astronomy Reports*, **56**, 589

Ikhsanov N.R. & Finger, M.H. 2012 *Astrophysical Journal*, **753**, 1

Correspondence principle is satisfied

$$|K_{\text{sd}}(r_m)| \leq \frac{\mu^2}{(r_m r_{\text{cor}})^{3/2}} \left(1 - \frac{\Omega(r_m)}{\omega_s} \right)$$

I. Accretion from a Turbulent-dominated Atmosphere

$\Omega(r_m) = 0$ $r_m \equiv r_A = \left(\frac{\mu^2}{\mathfrak{M} \sqrt{2GM_{\text{ns}}}} \right)^{2/7}$ $v_t(r_A) = v_k(r_A)$	$ K_{\text{sd}}^{(\text{t})}(r_A) \leq \mathfrak{M} \omega_s r_A^2$
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II. Accretion from a free-falling material (Bondi accretion scenarios)

$\Omega(r_m) = 0$ $r_m \equiv r_A$ $v_t(r_A) = \omega_s r_A$	$ K_{\text{sd}}^{(\text{ff})}(r_A) = K_{\text{sd}}^{(\text{t})}(r_A) \times \left(\frac{\omega_s r_A}{v_k} \right) \leq \frac{\mu^2}{r_{\text{cor}}^3}$
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Evaluation of the Magnetospheric Radius from Spin Evolution

$$|K_{\text{sd}}^{\max}| = \frac{k_t \mu^2}{(\textcolor{red}{r}_m r_{\text{cor}})^{3/2}} \geq 2\pi I |\dot{\nu}_{\text{sd}}|$$

$$\textcolor{red}{r}_m \leq \left(\frac{k_t \mu^2}{2\pi I |\dot{\nu}_{\text{sd}}|} \right)^{2/3} \left(\frac{\omega_s^2}{GM_{\text{ns}}} \right)^{1/3}$$

Name	B_{12} , G	P_s , s	$ \dot{\nu}_{\text{sd}} $, 10^{-12} Hz/s	r_m , 10^8 cm	r_A , 10^8 cm	r_m/r_A
OAO 1657-415	3.2	38	3.2	1.4	6.8	0.2
Vela X-1	2.6	283	0.3	1.5	8.3	0.18
4U 1907+09	2.1	438	0.04	2.9	6	0.48
GX 301-2	4	683	0.1	1.9	5.5	0.34
X Persei	3.3	837	0.024	4.5	23	0.2

Constraints on Accretion Scenario in LPXPs

Basic criteria:

- Slowly rotating accretion flow $\Omega(r_m) < \omega_s$
- Small magnetospheric radius $r_m \ll r_A$

I. Keplerian disk can be excluded

$$P_{\text{eq}} \sim 4 - 10 \text{ s} \ll P_{\text{obs}}$$

II. Quasi-spherical accretion can also be excluded

$$r_m^{(\text{sp})} \geq r_A$$

III. Magnetic accretion

- Accretion from a magnetized wind
- Formation of a non-Keplerian magnetized disk-like envelope (**magnetic slab**)

Some Historical Points of Magnetic Accretion Scenario

1959	Star formation in a magnetized medium	Mestel, L. 1959, MNRAS, 119, 223
1964 – 1965	Magnetic Collapse	Kardashev, N.S. 1965, Sov. Astron., 8, 643
1969	Fossil magnetic field in spherical flow onto a Neutron Star	Bisnovatyi-Kogan, G.S. Fridman, A.M. 1969, Sov. Astron. 13, 566
1971	Magnetic control of accretion onto a black hole flow deceleration at R_{sh}	Shvartsman, V.F. 1971, SvA, 15, 377
1974 – 1976	Non-Keplerian magnetic disk around a black hole	Bisnovatyi-Kogan, G.S., Ruzmaikin, A.A. 1974, ApSS, 28, 45 Bisnovatyi-Kogan, G.S., Ruzmaikin, A.A. 1976, ApSS, 42, 401
2003 – 2006	Non-Keplerian magnetic disk simulations for a Black Hole (magnetically arrested disk)	Igumenshchev, I.V., Narayan, R., Abramowicz, M.A. 2003, ApJ, 592, 1042 Igumenshchev, I.V. 2006, ApJ, 649, 361
2012 – 2013	Non-Keplerian magnetic disk around a Neutron Star (magnetic slab)	This talk (+ 7 papers by Ikhsanov, N.R. et al.)

Accretion from a magnetized flow $\left(\beta_0 = \frac{\mathcal{E}_{\text{th}}(R_{\text{G}})}{\mathcal{E}_{\text{m}}(R_{\text{G}})} \sim 1 \right)$

r-Ram $\mathcal{E}_{\text{ram}}(R_{\text{G}}) = \rho_{\infty} v_{\text{rel}}^2$

ϕ -Ram $\mathcal{E}_{\text{rot}}(R_{\text{G}}) = \rho_{\infty} (\Omega_{\text{orb}} R_{\text{G}})^2$

Thermal $\mathcal{E}_{\text{th}}(R_{\text{G}}) = \rho_{\infty} c_{\text{s}}^2(R_{\text{G}})$

Magnetic $\mathcal{E}_{\text{m}}(R_{\text{G}}) = \frac{B_{\text{f}}^2(R_{\text{G}})}{8\pi}$

$$\begin{aligned}\mathcal{E}_{\text{ram}}(\mathbf{r}) &= \mathcal{E}_{\text{ram}}(R_{\text{G}}) \left(\frac{R_{\text{G}}}{r} \right)^{5/2} \\ \mathcal{E}_{\text{rot}}(\mathbf{r}) &= \mathcal{E}_{\text{rot}}(R_{\text{G}}) \left(\frac{R_{\text{G}}}{r} \right)^{7/2} \\ \mathcal{E}_{\text{m}}(\mathbf{r}) &= \beta_0^{-1} \mathcal{E}_{\text{th}}(R_{\text{G}}) \left(\frac{R_{\text{G}}}{r} \right)^4\end{aligned}$$

Shvartsman radius

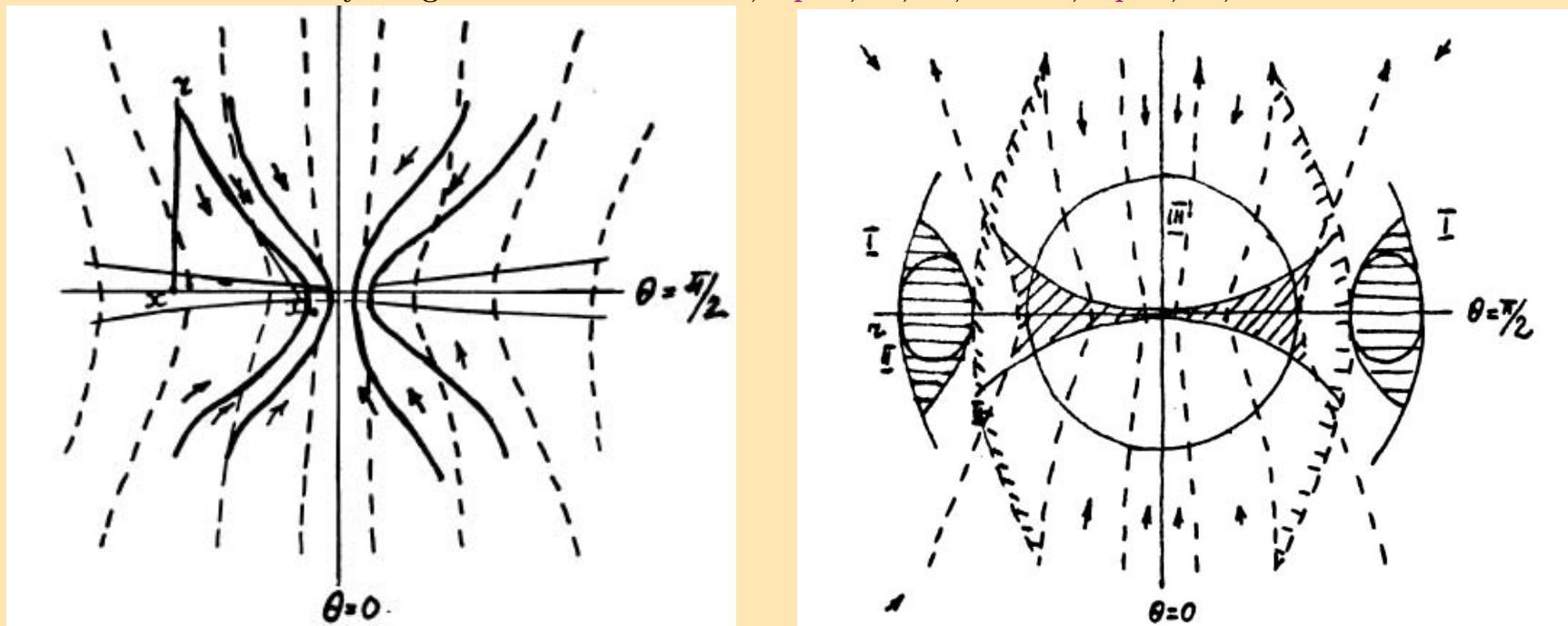
$$\mathcal{E}_{\text{m}}(R_{\text{sh}}) = \mathcal{E}_{\text{ram}}(R_{\text{sh}})$$

$$R_{\text{sh}} = \beta^{-2/3} \left(\frac{c_{\text{s}}}{v_{\text{rel}}} \right)^{4/3} R_{\text{G}}$$

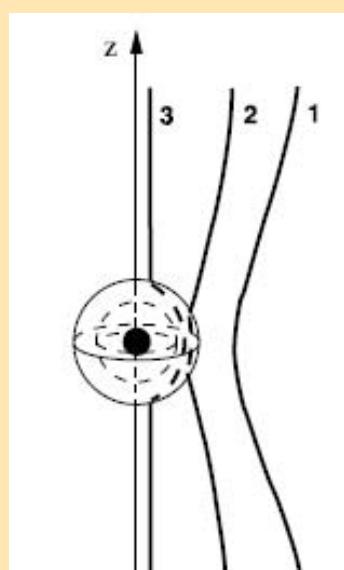
Non-Keplerian Magnetic Slab

$$t_{\text{rec}} = \frac{r}{\eta_m v_A} = \eta_m^{-1} t_{\text{ff}} \left(\frac{v_{\text{ff}}}{v_A} \right)$$

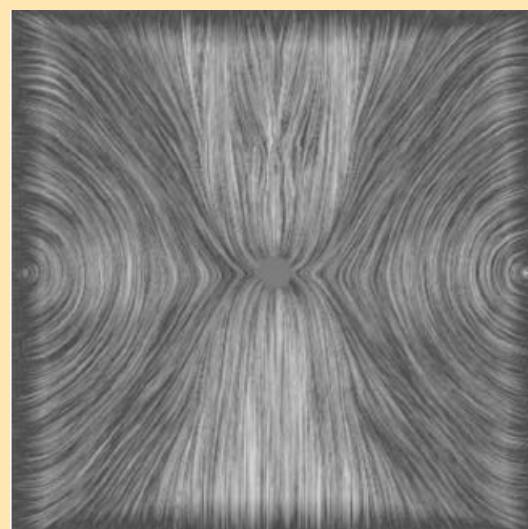
Bisnovatyi-Kogan & Ruzmaikin 1974, ApSS, 28, 45; 1976, ApSS, 42, 401



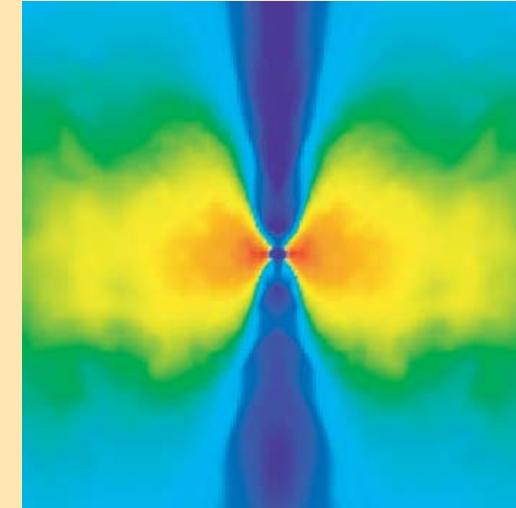
Igumenschev, Narayan & Abramowicz 2003, ApJ, 592, 1042



Magnetic field



Density



Magnetic accretion in X-ray pulsars

Basic condition:

$$R_{\text{sh}} > \max\{r_A, r_{\text{circ}}\} \rightarrow v_{\text{cr}} < v_{\text{rel}} < v_{\text{ma}}$$

$$v_{\text{ma}} \simeq 465 \text{ km s}^{-1} \times \beta_0^{-1/5} \mu_{30}^{-6/35} m^{12/35} \dot{m}_{15}^{3/35} \left(\frac{c_s(r_G)}{10 \text{ km s}^{-1}} \right)^{2/5}$$

$$v_{\text{cr}} \simeq 100 \text{ km s}^{-1} \times \beta_0^{1/7} \xi_{0.2}^{3/7} m^{3/7} P_{40}^{-3/7} \left(\frac{c_s(r_G)}{10 \text{ km s}^{-1}} \right)^{-2/7}$$

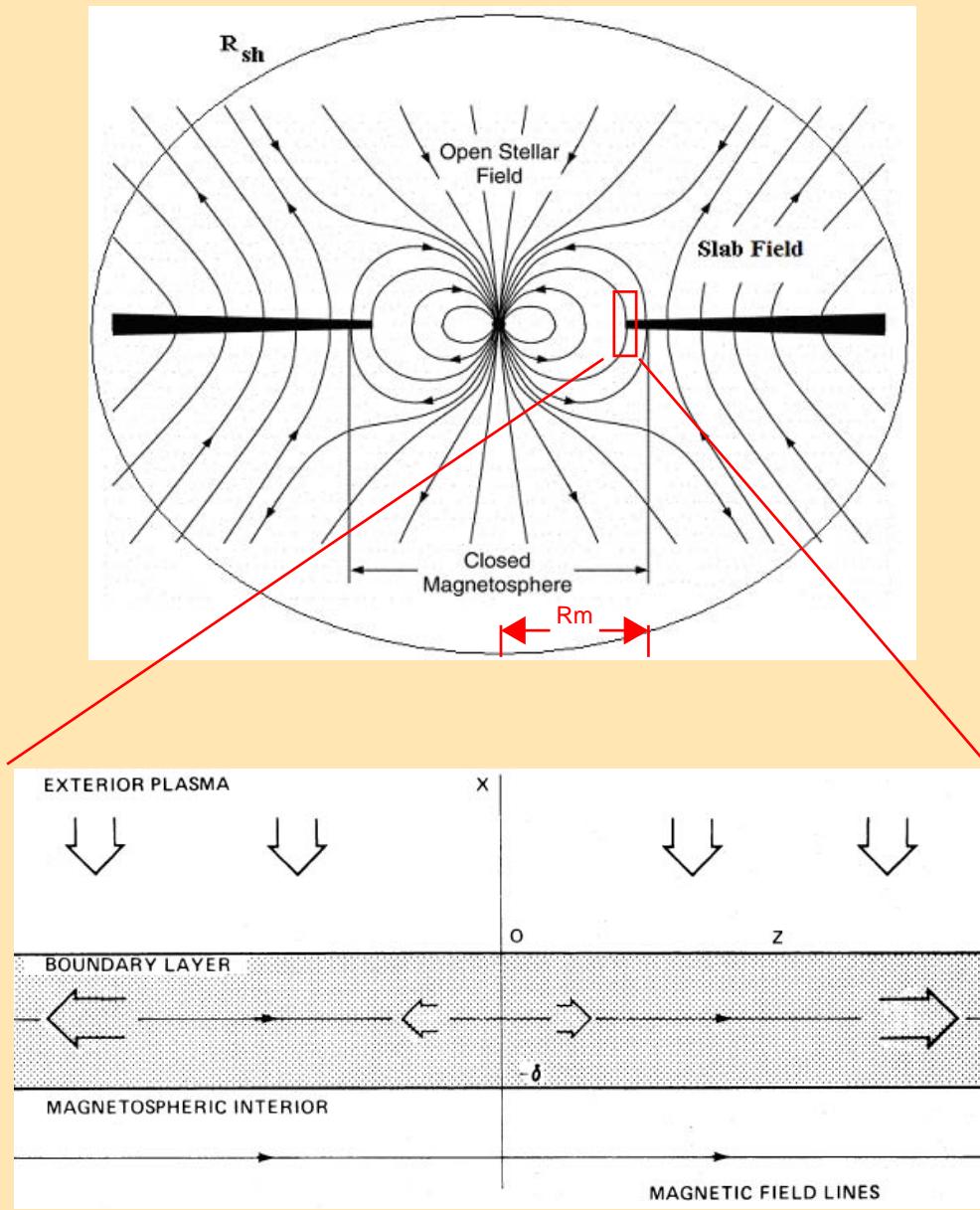
Three possible accretion scenarios

Quasi-spherical $\leftarrow v_{\text{rel}} > v_{\text{ma}}$

Keplerian disk $\leftarrow v_{\text{rel}} < v_{\text{cr}}$

Magnetic slab $\leftarrow v_{\text{cr}} < v_{\text{rel}} < v_{\text{ma}}$

Magnetic accretion onto a magnetized neutron star



$$c_s(r_m) > v_\phi(r_m)$$

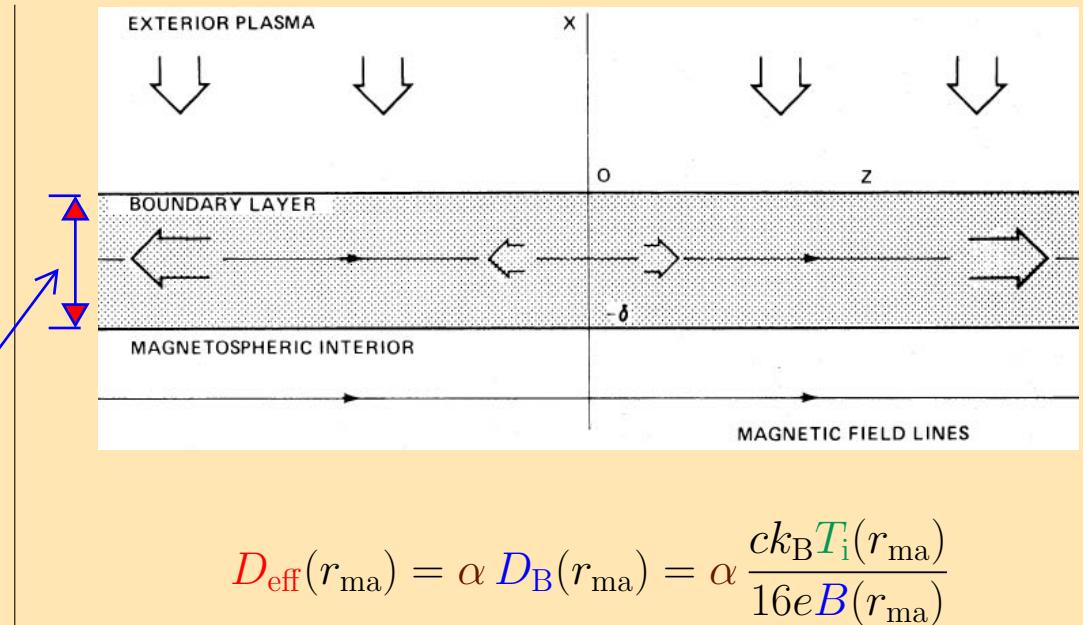
$$v_\phi(r_m) \ll v_k(r_m)$$

$$v_A(r_m) \sim c_s(r_m)$$

$$r_{ma} \leq r_m \leq r_A$$

Magnetospheric radius of a neutron star which accretes material from the non-Keplerian magnetic slab

$$\left\{ \begin{array}{l} \frac{\mu^2}{2\pi r_{\text{ma}}^6} = \rho(r_{\text{ma}}) c_s^2(r_{\text{ma}}) \\ \dot{\mathfrak{M}}_{\text{in}}(r_{\text{ma}}) = \frac{L_X R_{\text{ns}}}{GM_{\text{ns}}} \\ \dot{\mathfrak{M}}_{\text{in}}(r_{\text{ma}}) = 4\pi r_{\text{ma}} \delta_m \rho(r_{\text{ma}}) v_{\text{ff}}(r_{\text{ma}}) \\ \delta_m(r_{\text{ma}}) = \left[t_{\text{ff}}(r_{\text{ma}}) D_{\text{eff}}(r_{\text{ma}}) \right]^{1/2} \end{array} \right.$$



$$r_{\text{ma}} = \left(\frac{c m_p^2}{16 \sqrt{2} e k_B} \right)^{2/13} \frac{\alpha^{2/13} \mu^{6/13} (GM_{\text{ns}})^{5/13}}{T_0^{2/13} L_X^{4/13} R_{\text{ns}}^{4/13}}$$

Ikhsanov N.R. & Beskrovnyaya, N.G. 2012 *Astronomy Reports*, **56**, 589

Ikhsanov N.R., Kim, V.Y., Beskrovnyaya, N.G. & Pustil'nik, L.A. 2013 *Astrophysics & Space Sci.*, first on-line

Spin-down rates of LPXPs

$|\dot{\nu}_{\text{sd}}^{\text{max}}| \times 10^{-12} \text{ Hz s}^{-1}$

Name	B_{12} , G	P_s , s	$ \dot{\nu}_{\text{sd}}^{\text{obs}} $	$ \dot{\nu}_{\text{sd}}^{(0)} $	$ \dot{\nu}_{\text{sd}}^{(t)} $	$ \dot{\nu}_{\text{sd}}^{(m)} $
OAO 1657-415	3.2	38	3.2	0.056	0.18	3.3
Vela X-1	2.6	283	0.2	0.0003	0.012	0.4
4U 1907+09	2.1	438	0.04	0.0002	0.008	0.2
4U 1538-522	1.8	529	0.06	0.0002	0.008	0.15
GX 301-2	4	683	0.1	0.003	0.02	0.7
X Persei	3.3	837	0.024	0.0001	0.0013	0.03

$$|\dot{\nu}^{(0)}| = \frac{1}{2\pi I} \left[\frac{\mu^2}{r_{\text{cor}}^3} \right]; \quad |\dot{\nu}^{(t)}| = \frac{1}{2\pi I} \left[\mathfrak{M} \omega_s R_A \right]; \quad |\dot{\nu}^{(m)}| = \frac{1}{2\pi I} \left[\frac{\mu^2}{(r_{\text{ma}} r_{\text{cor}})^{3/2}} \right]$$

$$|\dot{\nu}_{\text{sd}}^{(0)}| < |\dot{\nu}_{\text{sd}}^{(t)}| \ll |\dot{\nu}_{\text{sd}}^{\text{obs}}| < |\dot{\nu}_{\text{sd}}^{(m)}|$$

Magnetic Accretion onto a Neutron Star

1. Accretion from a magnetized wind ($\beta_0 \sim 1$)
 2. Deceleration of the free-falling material at the Shvartsman radius R_{Sh}
 3. Formation of the non-Keplerian magnetic slab
 4. Accumulation and diffusion of material into the NS's magnetic field
 5. Stationary accretion at $\dot{\mathfrak{M}}_{\text{diff}}(r_m) = L_x R_{\text{ns}} / GM_{\text{ns}}$
-

New parameters:	Shvartsman radius $R_{\text{sh}} = \beta^{-2/3} \left(\frac{c_s(r_G)}{v_{\text{rel}}} \right)^{4/3} r_G$
	Magnetospheric radius $r_{\text{ma}} = \left(\frac{cm_p^2}{16\sqrt{2}e k_B} \right)^{2/13} \frac{\alpha_B^{2/13} \mu^{6/13} (GM_{\text{ns}})^{5/13}}{T_0^{2/13} L_x^{4/13} R_{\text{ns}}^{4/13}}$
	Spin-down torque $ K_{\text{sd}} = k_m \frac{\mu^2}{(r_{\text{ma}} r_{\text{cor}})^{3/2}} \left(1 - \frac{\Omega_{\text{sl}}(r_{\text{ma}})}{\omega_s} \right)$

Accretion-powered Be/X-ray Pulsar SXP 1062

Name	P_s , s	$\log L_x$	E_{cyc}	$\dot{\nu}$, Hz/s	P_{orb} , d	Sp. type	d
SXP 1062	1070	35.8	—	-2.6×10^{-12}	656	B0 IIIe	60 kpc

Magnetic field determination

(Ikhsanov N.R. 2012, MNRAS, 424, L39)

Spin-down torque	Magnetic field	Magnetospheric radius
$\frac{\mu^2}{r_{\text{cor}}^3} \geq 2\pi I \dot{\nu}_{\text{obs}}$	$B_* \geq 6 \times 10^{14} \text{ G}$	$r_A > r_{\text{cor}}$
$\mathfrak{M} \omega_s r_A^2 \geq 2\pi I \dot{\nu}_{\text{obs}}$	$B_* \geq 10^{15} \text{ G}$	$r_A > r_{\text{cor}}$
$\frac{\mu^2}{(r_{\text{ma}} r_{\text{cor}})^{3/2}} \geq 2\pi I \dot{\nu}_{\text{obs}}$	$B_* \geq 4 \times 10^{13} \text{ G}$	$r_{\text{ma}} \sim 0.01 r_{\text{cor}}$

Associated with a SNR of the age $\tau \sim (1 - 4) \times 10^4$ yr

Spin Evolution of SXP 1062

(Ikhsanov N.R. 2012, MNRAS, 424, L39)

Name	P_s , s	$\log L_x$	B_* 4×10^{13} G	$\dot{\nu}$, Hz/s	P_{orb} , d	Sp. type	d
SXP 1070	1062	35.8		-2.6×10^{-12}	656	B0 IIIe	60 kpc

Ejector phase <i>(Spin-powered pulsar)</i>	Magneto-dipole (conv.) Spitkovsky (2006) $\tau \sim 6 \times 10^4$ yr	
	Beskin et al. (1993) $\tau \sim 10^4$ yr	
		$\tau \sim 10^4$ yr
Propeller phase	$\tau_{\text{prop}} \sim \frac{\pi I r_{\text{ma}}^3}{\mu^2 P_{\text{ej}}}$	$\tau_{\text{prop}} \sim 110$ yr
Accretor phase	$\tau_{\text{acc}} \sim \frac{1}{2 \dot{\nu}_{\text{obs}} P_{\text{prop}}}$	$\tau_{\text{acc}} \sim 600$ yr
Age of the Pulsar is $\tau \sim (1 - 2) \times 10^4$ yr		

Young Isolated X-ray Pulsar 1E 161348-5055

Name	P_s , s	L_x , erg/s	E_{cyc}	$ \dot{\nu} $, Hz/s	Type	d, kpc
1E 1613	24030	10^{34}	—	3×10^{-18}	Isolated	3.3

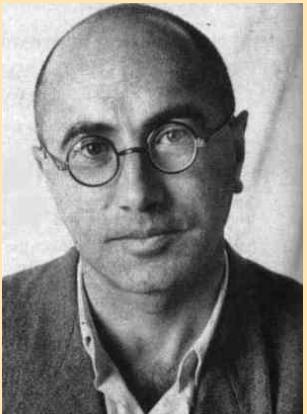
Magnetic field determination (Ikhsanov N.R., et al. 2013, ApSS, 346, 105)

Spin-down scenario	Magnetic field	Magnetospheric radius
Ejector (spin-powered pulsar)	$B_* \geq 4 \times 10^{18} \text{ G}$	$r_A > 100 r_{\text{cor}}$
Keplerian fossil disk	$B_* \geq 10^{16} \text{ G}$	$r_A > 2.5 r_{\text{cor}}$
Magnetic fossil slab	$B_* \geq 8 \times 10^{11} \text{ G}$	$r_{\text{ma}} \sim 0.002 r_{\text{cor}}$

$$B \geq 10^{12} \text{ G} \left(\frac{\tau}{2000 \text{ yr}} \right)^{-13/17} \times k_m^{-13/17} \alpha_B^{3/17} I_{45}^{13/17} m^{8/17} T_6^{3/17} \mathfrak{M}_{14}^{-6/17}$$

Магнитная реверсная аккреция: параметры источников

Название	P_s , с	$ \dot{\nu}_{sd} $, $10^{-12} \text{ Гц с}^{-1}$	L_X , $10^{34} \text{ эрг с}^{-1}$	B_* , 10^{12} Гс	r_{cor} , 10^8 см	r_{ma} , 10^8 см	a_p , 10^5 см	T_{bb} , кэВ
SGR 1627-41	2.6	2.8	0.25	3.9	3.2	1.2	0.9	0.5
SGR 1900+14	5.2	3.4	9.0	2.7	5.2	1.0	1.0	1.1
SGR 1806-20	7.6	13	16	12	6.6	1.7	0.8	1.4
SGR 0526-66	8.05	0.59	14	1.2	6.9	0.6	1.3	1.1
1E 1547.0-5408	2.07	11	0.08	0.63	2.8	2.1	0.7	0.4
CXOU J174505.7-381031	3.83	4.4	6.0	2.3	4.2	1.0	1.0	0.98
PSR J1622-4950	4.33	0.9	0.063	1.5	4.6	1.2	0.9	0.33
XTE J1810-197	5.54	0.26	3.9	0.3	5.4	0.46	1.5	0.73
1E 1048.1-5937	6.45	0.55	0.6	0.31	5.95	0.83	1.1	0.53
1E 2259+586	6.98	0.01	2.2	0.024	6.3	0.17	2.4	0.49
CXOU J010043.1-721134	8.02	0.29	6.1	0.51	6.9	0.51	1.4	0.83
4U 0142+61	8.69	0.03	11	0.11	7.26	0.21	2.2	0.77
CXO J164710.2-455216	10.61	0.006	0.3	0.011	8.3	0.22	2.1	0.32
1RXS J170849.0-400910	11.0	0.16	5.9	0.4	8.5	0.46	1.5	0.81
1E J1841-045	11.8	0.28	19	0.99	8.9	0.49	1.4	1.1



Yakov Zel'dovich



Victory Shvartsman

Final Remarks

A Farewell to Magnetars...

Welcome to Magnetic Accretion



Gennady Bisnovatyi-Kogan



Alex Ruszmaikin